

Haptic Perception of Viscoelastic Materials

by

Ozan aldıran

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Doctor of Philosophy

in

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Haptic Perception of Viscoelastic Materials

Koç University

Graduate School of Sciences and Engineering

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Ozan Çaldıran

and have found that it is complete and satisfactory in all respects,
and that any and all revisions required by the final
examining committee have been made.

Committee Members:

Prof. CAGATAY BASDOGAN (Advisor)

Prof. HONG Z. TAN (Co-Advisor)

Assoc. Prof. MEHMET SAYAR

Assoc. Prof. WOUTER M. BERGMANN TIEST

Prof. BURAK GÜÇLÜ

Assoc. Prof. KNUT DREWING

Date: _____



To my beloved family...

ABSTRACT

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In our daily lives, we interact with different types of deformable materials. Regarding their mechanical behavior, some of those materials lie in a range that is between purely elastic and purely viscous. This range of mechanical behavior is described as viscoelasticity. In certain types of haptic interactions such as assessment of ripeness of fruit, firmness of cheese, and consistency of organ tissue, we rely heavily on our haptic perception of viscoelastic materials. The relationship between the mechanical behavior of viscoelastic materials and our perception of them has been investigated in the field of psychorheology. However, our knowledge on how we perceive viscoelastic materials is still quite limited though some research work has already been done on purely elastic and purely viscous materials. History- and frequency-dependent behavior of viscoelastic materials results in a complex time-dependent response, which requires relatively more sophisticated models to investigate their behavior than those of purely elastic and viscous materials. In this study, we model viscoelasticity using a “springpot” (i.e., fractional order derivative element) and express its behavior in the frequency domain using two physical parameters: “magnitude” and “phase” of complex stiffness. In the frequency domain, we are able to devise signal detection experiments where we can investigate the perception of viscoelastic materials using the perceptual terms of “firmness” and “bounciness”, corresponding to the physical parameters of “magnitude” and “phase”. The results of our experiments show that the JND for bounciness increases linearly with increasing “phase”, following Weber’s law, while the JND for firmness is surprisingly independent of the level of “phase”.

ÖZETÇE

Viskoelastik Materyallerin Dokunsal Algısı

Ozan Çaldıran

Makine Mühendisliği, Doktora

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Günlük hayatımızda değişik tipte deformasyon özellikleri gösteren maddeler ile etkileşime geçmekteyiz. Bu maddelerin bazıları mekanik davranışları bakımından saf elastikiyet ile saf vizkozite arasında bir yelpazade yer almaktadır. Bu mekanik davranış yelpazesi viskoelastisite olarak tanımlanmaktadır. Bir meyvenin olgunluğunu, peynirin sıklığını veya iç organların sertliğini anlamaya çalışma gibi belirli dokunsal etkileşim durumlarında dokunma duyumuzun viskoelastik materyalleri algılama kapasitesine dayanmaktayız. Psikoreoloji bilimdalı viskoelastik materyallerin mekanik davranışları ile onları algılayış şeklimiz arasındaki ilişkiyi incelemektedir. Fakat, viskoelastik materyalleri nasıl algıladığımıza dair bilimiz hala oldukça kısıtlı. Viskoelastik materyallerin deformasyon geçmişi ve uyarılma frekansına dayalı olarak sergilediği mekanik davranışlar oldukça karmaşık sonuçlar ortaya çıkarabilmektedir ve araştırılmaları ancak gelişmiş modeller aracılığıyla mümkün olmaktadır. Bu çalışmada viskoelastisiteyi kesir dereceli türev elemanı ile modelledik ve bu elemanın davranışını frekans uzayında ifade ederek iki fiziksel değişken (genlik ve faz) cinsinden tanımladık. Bu sayede, genlik değişkeni için "sıklık" ve faz değişkeni için ise "geri sekerlik" terimlerini kullanarak viskoelastik materyal algısını sinyal tespit teorisi deneyleri gerçekleştirerek inceleyebildik. Deney sonuçları geri sekerlik için ancak farkedilebilir farkın Weber kuralı ile uyumlu şekilde faz ile doğrusal olarak arttığını gördük. Öte yandan sıklık değişkeninin ancak farkedilebilir farkı ise değişik faz seviyeleri için farklılık göstermedi.

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ABBREVIATIONS

SDT	Signal Detection Theory
JND	Just-noticeable difference
WF	Weber fraction
CR	Correct reject
FA	False alarm
M	Miss
H	Hit
GLD	Grunwald–Letnikov derivative
STM	Short-term memory
DFT	Discrete Fourier Transform
FOP	Force over Position

Chapter 1

INTRODUCTION

Viscoelastic material behavior is commonly observed in nature especially in biological materials [1]. In interaction scenarios such as squeezing a piece of fruit to assess its ripeness or palpating soft organ tissues to assess their consistency, we have to rely on our haptic perception of viscoelasticity. In certain settings, our judgments about the viscoelastic properties of these materials can be of critical importance. For example, tissue consistency is one of the four criteria (the 4 Cs) used by surgeons to assess muscle debridement alongside color, contractility, and capacity to bleed; though research shows that decisions made based on the 4 Cs do not agree with histopathological findings [2]. Investigating the discrimination thresholds relating to our haptic perception of material consistency will help us to gain insight into the reliability of our judgments in such critical settings.

Besides, the design of medical simulation and teleoperation systems can also benefit from the knowledge of the perceptual thresholds. Interactive medical simulations require real-time computation of force response of a modeled tissue. However, developers may not be able to fulfill this requirement unless they make a sacrifice in the fidelity of the models [3]. In this regard, our perceptual thresholds can be used as a guideline for determining this fidelity. Similarly, teleoperation systems require real-time transmission of data between local and remote sites. One approach to reduce data transfer and achieve real-time interaction is to exploit the limitations of the human perceptual system. Researchers have proposed data reduction techniques based on the just-noticeable differences (JNDs) relevant to haptic modality [4, 5]. These techniques can be tailored to the medical domain, or any other domain relating to viscoelastic materials, using a systematic knowledge on the thresholds of

haptic perception of viscoelasticity.

However, most studies on haptic perception of deformable materials have focused on either elasticity or viscosity, but not both [6]. Researchers scrutinized the question of how humans judge the stiffness of a spring or the viscosity of a liquid. They have investigated the prevalent haptic cues used in stiffness or viscosity judgment [7, 8, 9, 10, 11]; integration of information coming from different cues or senses [12, 8]; and effects of time delay in force feedback [13, 14]; exploration strategies [15, 16], tool use [15, 17, 9]; boundary crossings [10, 11]; the relation of perception and action, and effects of arm posture and force direction [18, 19].

Compared to viscoelasticity, both elasticity and viscosity are relatively straightforward to investigate since they are governed by constant ratios of the force response to displacement (elasticity) or to velocity (viscosity). On the other hand, viscoelasticity is a highly complex, history- and frequency-dependent material behavior. Viscoelastic materials can show both energy storage and loss characteristics depending on the loading frequency and history. Stress-relaxation and creep are two phenomena where we observe these two characteristics. Moreover, a phase difference between force and displacement occurs for a viscoelastic material under dynamic loading. In the next section, we will explain these characteristics of viscoelastic materials.

1.1 Background on Viscoelasticity

1.1.1 Stress-Relaxation

If a viscoelastic material is compressed at a sufficiently high rate and then held at the final displacement (ramp and hold input), we can observe a relaxation of the force response. Initially, a portion of the strain energy is stored, similar to an elastic material, but then some portion of this energy is dissipated (see Fig. 1.1a). In a quasi-static case, where the compression rate is very low, the dissipation will occur during the compression, and no relaxation will be observed (see the dashed lines in Fig. 1.1a).

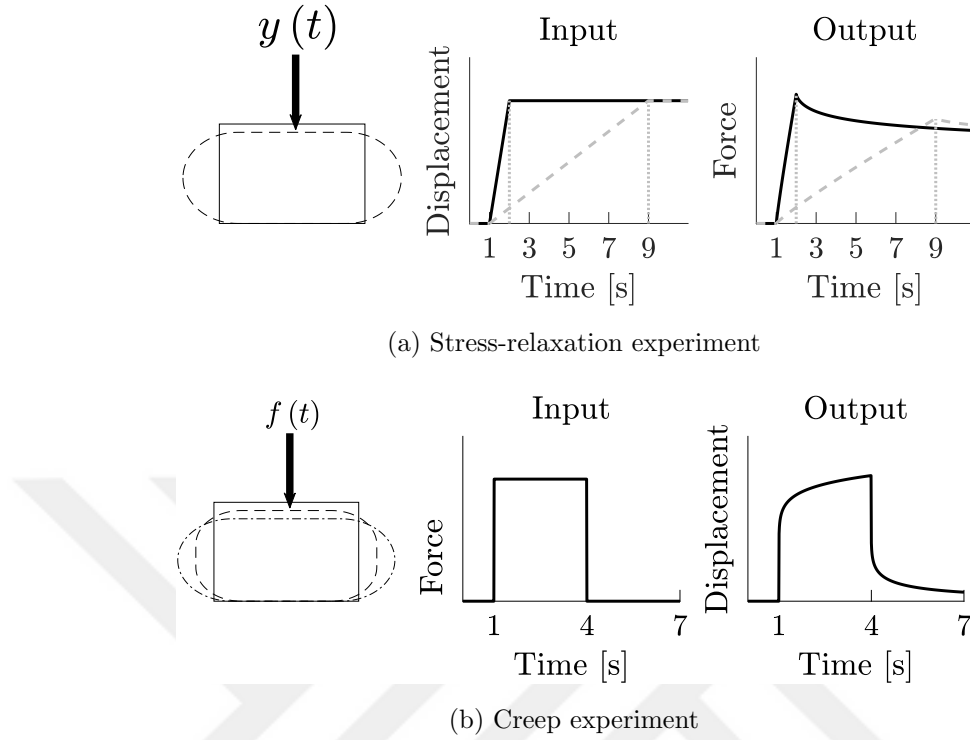


Figure 1.1: Mechanical behavior of viscoelastic materials. (a) In stress-relaxation experiments, ramp and hold displacement inputs at various rates (see the middle panel) are applied to the tested materials (see the left panel, where the undeformed and deformed states of a specimen are drawn with solid and dashed lines, respectively). The force responses are analyzed for the magnitude at steady-state and the decay rate (see the right panel). In the middle and right panels, the solid black and dashed grey curves represent the cases for a high and low compression rate, respectively. (b) In creep experiments, step force inputs (see the middle panel) are applied to the tested materials. In the left panel, the undeformed and deformed states of a specimen are drawn with solid, dashed (onset of input) and dash-dotted (offset of input) lines, respectively. The restoration of the material to its undeformed state is observed (see the right panel).

1.1.2 Creep

Creep can be seen as the dual of stress-relaxation. In a creep experiment, the controlled input is a step force applied for a certain amount of time and released afterward. Initially, we observe an instantaneous deformation like an elastic material. As time progresses, the material further deforms but with an exponentially decaying strain rate. At the instant when the force is released, the initial portion of the strain is recovered, but full recovery takes more time (see Fig. 1.1b).

1.1.3 Phase Difference

A phase difference occurs between the force and displacement of viscoelastic materials stimulated under a dynamic load (see Fig. 1.2). The magnitude of this difference can be used to determine the relationship between the energy loss and storage characteristics of a viscoelastic material. The smaller (larger) the phase difference, the more restorative (dissipative) the material is.

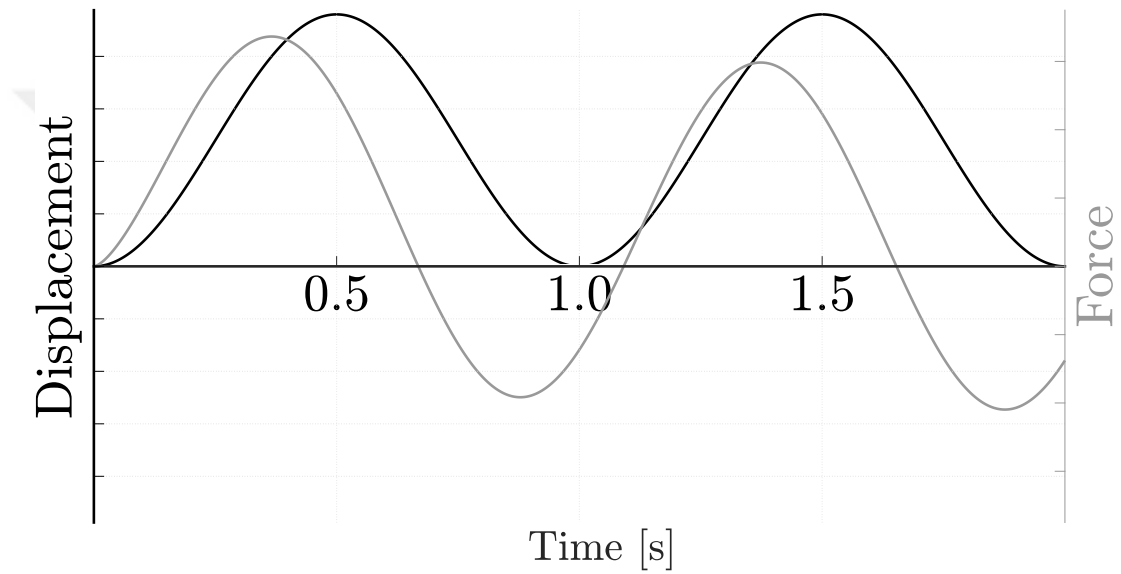


Figure 1.2: Phase difference between the displacement (the black curve) and force response (the grey curve) of a viscoelastic material.

1.2 Literature Review

Despite the common occurrence of viscoelastic materials in nature, only a few studies focused on the haptic perception of viscoelasticity. We believe that this is partly due to the increased complexity of the models required to accommodate the above described characteristics. The increase in complexity is reflected in the number of model parameters and the coupled effect of these model parameters on the mechanical response of a viscoelastic material complicates also psychophysical investigation. For example, Nicholson et al. [20] used an electromechanical device mimicking the Kelvin-Voigt model, a spring and a dashpot connected in parallel, to investigate

whether palpation is a reliable diagnostic tool to examine pathologies of human spine. Although Kelvin-Voigt model involves two coefficients (one stiffness and one damping coefficients), the authors used only the stiffness component of the model in the analysis of the experimental data. Similarly, Fakhoury et al. used a Maxwell arm, a spring and dashpot in series, for the characterization of their rubber stimuli, but once again only the stiffness component of the model was considered for interpreting the results of their psychophysical experiments [17].

Studies focused on the effect of delay on perception of stiffness are also of particular relevance to this study since viscoelastic materials show a phase difference between force and displacement under dynamic loading. In [21], it was shown that stiffness is underestimated when force is leading the displacement, and overestimated when force is delayed. Later in [13], several cues were proposed as possible candidates to explain this shift in perception. Among them, the ratio of peak force to penetration depth was found to be the most likely candidate. The study was further extended in [22] to model the motor control under delay, and the authors defined motormetric curves analogous to psychometric curves. In [23], these motormetrics were used to assess the discrepancies between perception and action during a teleoperated needle insertion task, where soft tissue was modeled as a spring with time-delay in the response, and a nonlinear boundary contact layer. Nonetheless, such model simplifications result in a loss of fidelity in modeling and is not suitable for our purposes.

In a preliminary work, we have proposed to investigate viscoelastic materials in the frequency domain for psychophysical evaluation without any sacrifice in modeling [24]. Such a representation enables us to alter only a single physical quantity, magnitude or phase difference through the back transformation to the parameter space, irrespective of the chosen viscoelastic model (see Fig. 1.2). Utilizing this approach, we investigated points of subjective equality between a Maxwell arm and purely elastic and viscous models. Similar approaches have also been utilized before for investigating the haptic perception of dynamic systems comprised of mass, spring, and damper [14, 25, 26, 27]. Although these approaches require participants

to follow a prescribed sinusoidal motion, they allow studying dynamic behavior that is otherwise not achieved so far according to the best of our knowledge. Moreover, such sinusoidal movement patterns can be viewed as application and release of pressure applied to deformable materials for assessing their softness [28], and repeating such loading-unloading cycles increases the sensory information gathered this way [29, 30, 31, 32, 33].

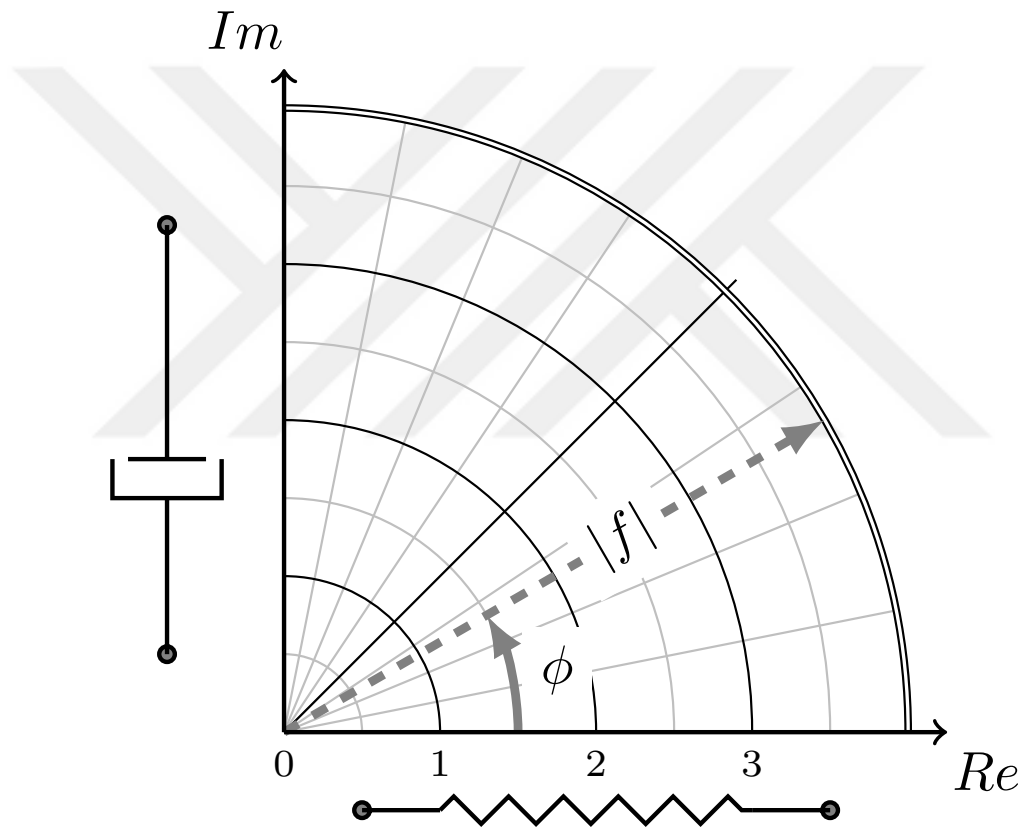


Figure 1.3: Viscoelastic spectrum in polar coordinates. The angular coordinate describes the phase difference ϕ between force f and displacement, while the radial coordinate describes the magnitude of the force response $|f|$. A spring (damper) can be equivalently represented by a viscoelastic material with purely real (imaginary) components.

Although our earlier approach [24] made the psychophysical evaluation of viscoelastic material behavior possible, we have not utilized any perceptual terms to assess the subjective feelings of participants. Fortunately, in psychorheology, a branch

of psychophysics that specializes in materials that flow such as viscous and viscoelastic materials, this issue has already been addressed [34]. In the psychorheology field, researchers have proposed lists of adjectives to describe the perceptual properties of rheological materials [35, 36, 37, 38]. They also tried to establish links between the physical aspects of the models and the adjectives through descriptive analysis [38, 39]. These links were then used to measure quality, acceptance or discriminability of certain products concerning the haptic sensory experience of consumers, especially in food and cosmetics research [40]. For example, researchers studied the effects of the fat content of dairy products [41], aging of cheese [42], or different brands of body lotions [43] on our perception.

In one of the earliest works of the psychorheology field, Scott-Blair and Coppen defined the “firmness” for a ball of California bitumen, which is a viscous material, as well as an elastic rubber specimen using a unified formulation [44]. The term “firmness” was chosen over “hardness” because the latter has different meanings in other branches of science [35, 36, 37]. Moreover, Scott-Blair and Coppen’s realizations about the temporal characteristics of viscoelastic materials also led to the introduction of a new rheological model called springpot (Scott-Blair) element, which is described by fractional order derivatives [45, 46, 47].

1.2.1 Perception of Purely Elastic and Purely Viscous Materials

Under the light of the discussion above, we can view viscoelasticity as a spectrum of material behavior spanned from pure elasticity to pure viscosity (see Fig. 1.2). Hence, investigating JND of viscoelastic materials can provide us with insights about how JND varies along this spectrum since the earlier studies focused on its lower (elasticity) and upper (viscosity) bounds only. Tan et al. measured the compliance (reciprocal of stiffness) JND with a signal detection experiment [48], whereas Beauregard used the same method to determine viscosity JND [49]. Also, Jones and Hunter investigated both stiffness and viscosity JNDs with contralateral limb-matching experiments [50, 51]. In these studies, the Weber fractions (WFs) reported for viscosity (13.6 % in [49] and 34 % in [51]) were higher than the WFs for elasticity

(8 % in [48] and 23 % in [50]).

Perception studies on elasticity and viscosity have also determined the respective JNDs under more diverse conditions. In Table 1.1, JNDs of elastic and viscous materials are presented as Weber fractions (WF)¹. However, it should be noted that this table includes only includes some exemplar studies, and a vast literature exists on the JND of elastic and viscous materials.

In the present study, we left the confines of the viscoelastic spectrum and probed further into it regarding the discrimination thresholds of “bounciness” and “firmness”. Both terms were adopted from the psychorheology field. The terms “bounciness” and “firmness” describe the perceived phase difference and magnitude of a viscoelastic material response under displacement, respectively. In the next section, more detailed descriptions of these terms are given after mathematically describing the viscoelastic spectrum in terms of a integer order model, Maxwell model, and a fractional order model, springpot element. These descriptions inform our decision of the use of springpot element to study the haptic perception of viscoelastic materials.

For psychophysical evaluations, we designed two signal detection experiments and measured the JNDs of perceived “firmness” and “bounciness” of the springpot element. In the experimental design, we used the frequency-domain approach introduced in our earlier study [24]. In the bounciness experiment, in line with earlier studies on the perception of dynamic systems, we visually displayed reference movement trajectories with a constant frequency. Thus, visual feedback was always present in our experiments. In that sense, our study falls in the category of visuo-haptic perception. Also, the stimuli were haptically rendered by a point-contact electromechanical device which allows us to tune the viscoelastic model parameters freely at the expense of reduced tactile information.

¹Weber fraction is the ratio of JND to the reference stimulus intensity. Although Weber fraction is a convenient way of comparing JNDs, it may not always hold [52].

Table 1.1: Elasticity and Viscosity JND as Weber Fractions (WF)

Effect of		Publication	Stimulus	Feedback	WF
Visual feedback		[53]	Elastic	P	6–8 %
				P + V	5 %
		[54]	Elastic	P	14.2 %
				P + V	17.2 %
		[55]	Elastic (Hard)	H	18 %
				H + V	23 %
				V	200 %
			Elastic (Soft)	H	16 %
				H + V	15 %
				V	38 %
Cutaneous feedback	[8]	Elastic	H	15 %	
			P	50 %	
Material type		[48]	Elastic	P	8 %
		[49]	Viscous	P	13.6 %
		[50]	Elastic	P	23 %
		[51]	Viscous	P	34 %
Displacement	Fixed	[48]	Elastic	P	8 %
	Roving			P	22 %
Feedback:		P: Proprioceptive, V: Visual H: Haptic (Proprioceptive and Tactile)			

Chapter 2

VISCOELASTIC MODELS IN THE FREQUENCY DOMAIN

We use rheological models to mathematically represent the mechanical behavior of viscoelastic materials. Rheological models are specific configurations of spring, dashpot, and springpot elements. Integer derivative orders of 0 and 1 describe the force-displacement relationships of spring and dashpot (damper) elements, respectively. However, we require fractional order derivatives to model the same relationship in the case of springpot elements. Despite the computational complexity of the fractional order differentiation, rheological models utilizing a springpot element introduces several advantages over the ones composed of springs and dashpots alone (see applications in tissue modeling in [56, 57, 58], for example).

We analyze both an integer-order, a Maxwell arm, and a fractional-order rheological model, a springpot element, in the frequency domain to highlight the advantages of the use of fractional models for psychophysical investigation.

2.1 Maxwell Arm

Generalized Maxwell Model (GMM) (see Fig. 2.1) is a proper viscoelastic model which can capture time and rate dependency of viscoelastic materials. One of the foundational units of a GMM is a Maxwell arm: a spring and a damper in series. Each Maxwell arm of a GMM represents a different time constant τ of the modeled material. In this section, we investigate the mechanical response of a Maxwell arm with a large or a small time constant, i.e., the complex stiffness is close to being purely real or imaginary.

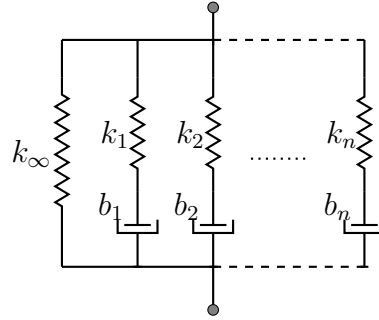


Figure 2.1: Generalized Maxwell Model with n arms is typically utilized to model viscoelastic materials. Each parallel unit composed of a serially connected spring and a damper is called a Maxwell arm.

The governing equation of a Maxwell arm is given as

$$\frac{\dot{f}}{k_1} + \frac{f}{b_1} = \dot{y} \quad (2.1)$$

where y is the displacement applied to the arm ($\dot{y}$ represents its time derivative); f is the force applied to the arm, $\dot{f}$ is the time derivative of the force; and k_1 and b_1 are the stiffness and damping coefficients of the spring and the damper, respectively.

We can transform (2.1) into Laplace domain and re-arrange it to obtain the transfer function between the displacement $Y(s)$ and force $F(s)$:

$$\frac{F(s)}{Y(s)} = \frac{\tau k_1 s}{\tau s + 1} \quad (2.2)$$

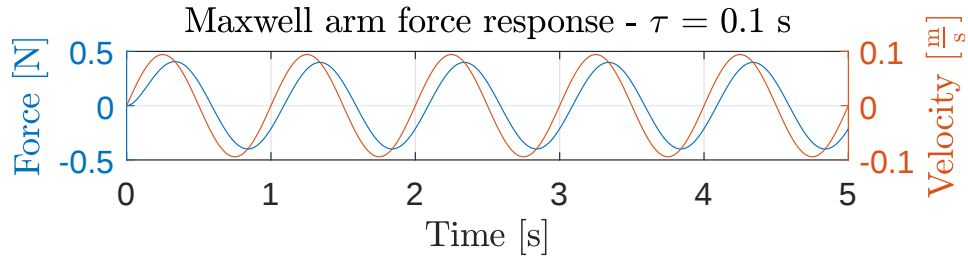
where $\tau = b_1/k_1$.

2.1.1 Limit Cases

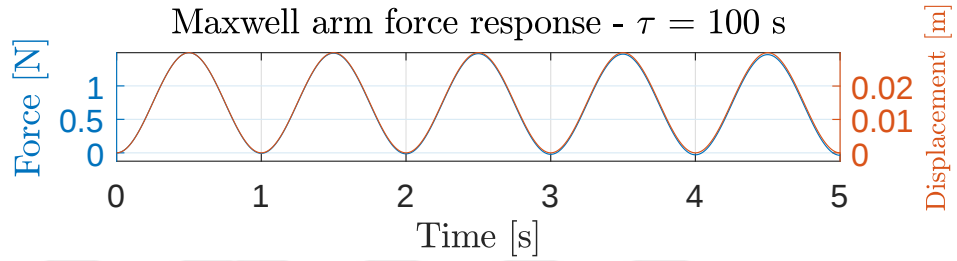
If we look at the transfer function in (2.2) when $\tau \rightarrow 0$:

$$\begin{aligned} \lim_{\tau \rightarrow 0} F(s) &= \lim_{\tau \rightarrow 0} \frac{\tau k_1 s}{\tau s + 1} Y(s) = \lim_{\tau \rightarrow 0} \frac{b_1/k_1 k_1 s}{\tau s + 1} Y(s) \\ &= b_1 s Y(s), \end{aligned} \quad (2.3)$$

we can see that a Maxwell arm behaves like a damper (see Fig. 2.2a).



(a) Viscous behavior



(b) Elastic behavior

Figure 2.2: Force response of a Maxwell arm with (a) small and (b) large time constants.

On the other hand, when $\tau \rightarrow \infty$, the response of a Maxwell arm is equivalent to that of a spring (see (2.4) and Fig. 2.2b).

$$\begin{aligned}
 \lim_{\tau \rightarrow \infty} F(s) &= \lim_{\tau \rightarrow \infty} \frac{\tau k_1 s}{\tau s + 1} Y(s) \\
 &= \lim_{\tau \rightarrow \infty} \frac{k_1 s}{s + 1/\tau} Y(s) \\
 &= k_1 Y(s)
 \end{aligned} \tag{2.4}$$

2.1.2 Harmonic Response of a Maxwell Arm

To look at the harmonic response of a Maxwell arm, we can plug-in $j\omega$ for s in the transfer function given in (2.2) to obtain the complex stiffness as

$$k_C(\omega) = \frac{\tau k_1 j \omega}{1 + \tau j \omega}, \quad (2.5)$$

$$k_C(\omega) = \underbrace{\frac{k_1 \omega^2 \tau^2}{\omega^2 \tau^2 + 1}}_{k_S} + j \underbrace{\frac{k_1 \omega \tau}{\omega^2 \tau^2 + 1}}_{k_L}. \quad (2.6)$$

The above formulation of complex stiffness is similar to the one given in (2.6), and can be alternatively written in phasor notation as

$$\begin{aligned} k_C(\omega) &= |k_C| \angle \phi \\ \phi &= \tan^{-1}(k_L/k_S) = \tan^{-1}(1/\omega\tau) \\ |k_C| &= \sqrt{k_S^2 + k_L^2}. \end{aligned} \quad (2.7)$$

In Fig. 2.3, a contour plot of the phase angle for a range of frequency ω and time constant τ pairs is shown. The yellow regions below the black curve correspond to a more damper-like behavior while the red regions above the blue curve correspond to a more spring-like behavior. The behavior in the orange region (in between the black and blue curves) is neither damper nor spring-like.

2.2 Springpot Element

Springpot element is the basic building block of fractional viscoelastic models [59]. These models can capture diverse complex material behavior. Regardless, a springpot element can also serve as a representative model reflecting the common viscoelastic material characteristics: stress-relaxation, creep, and phase difference under dynamic loading (see Fig. 1.1 and 2.4).

For a springpot element, the relationship between force f and displacement y is formulated as

$$f = \chi D_t^\alpha y, \quad (2.8)$$

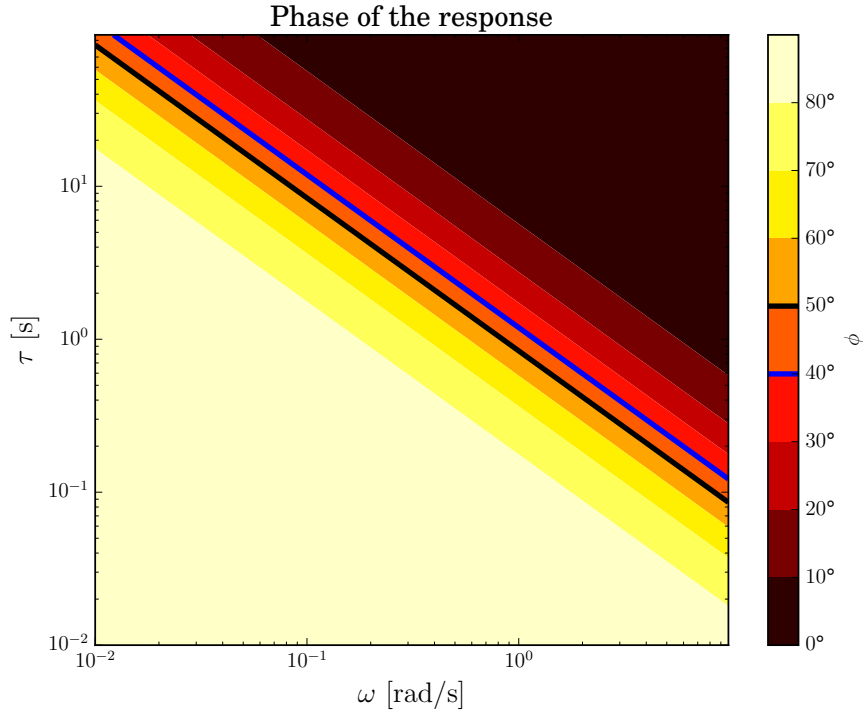


Figure 2.3: Contour plot of the phase angle ϕ with respect to the time constant τ vs frequency ω in log scale. The behavior is damper-like in regions under the black curve (shades of yellow), spring-like in regions above the blue curve (shades of red), and neither damper nor spring-like in the region between the black and blue curves (orange colored region).

where $D_t^\alpha = (d^\alpha)/(dt^\alpha)$ is the fractional derivative operator of order α with respect to time t , and χ is a quasi-property or coefficient whose units change with the derivative order.

Fractional derivatives can be thought of as a generalization of the more familiar concept of integer order derivatives to fractional number orders. For example, the fractional derivative of order one with respect to time applied to displacement y ($D_t^1 y$) is still the first derivative of y , i.e., velocity $\dot{y}$. So, a springpot element with $\alpha = 0$ corresponds to a spring, and $\alpha = 1$ corresponds to a damper. Moreover, if we analyze the force response of a springpot element in the frequency domain, we can almost decouple the contributions of the variables χ and α [60]. For a one-dimensional springpot element under a dynamic load of $y(t) = Y e^{i\omega t}$ with

magnitude Y and frequency ω , Eq. (2.8) becomes

$$f(t) = \chi \frac{d^\alpha Y e^{i\omega t}}{dt^\alpha}, \quad (2.9)$$

where $f(t)$ is the force response, and i represents a complex number. Differentiation of the displacement $y(t)$ leads to

$$f(t) = \chi (i\omega)^\alpha Y e^{i\omega t}, \quad (2.10)$$

Here, one can replace i^α with its Euler form:

$$f(t) = \chi \omega^\alpha Y (e^{i\pi/2})^\alpha e^{i\omega t}. \quad (2.11)$$

or more compactly,

$$f(t) = \chi \omega^\alpha Y e^{i(\omega t + \phi)}, \quad (2.12)$$

where the phase difference $\phi = \alpha \pi/2$. In Fig. 2.4, force responses of springpot elements with varying derivative orders α are drawn with respect to the same displacement input.

We can write Eq. (2.12) in terms of the complex stiffness k_C as a function of ω :

$$f(t) = k_C(\omega) Y e^{i\omega t}. \quad (2.13)$$

So, the relationship between the force response and displacement input can be represented by a magnitude $|k_C| = \chi \omega^\alpha$ and phase difference ϕ . Consequently, we can represent the viscoelastic spectrum of a springpot element in polar coordinates (see Fig. 2.5). This spectrum is defined from 0 to 1 in terms of α . For $\alpha = 0$, Eq. (2.12) reduces to

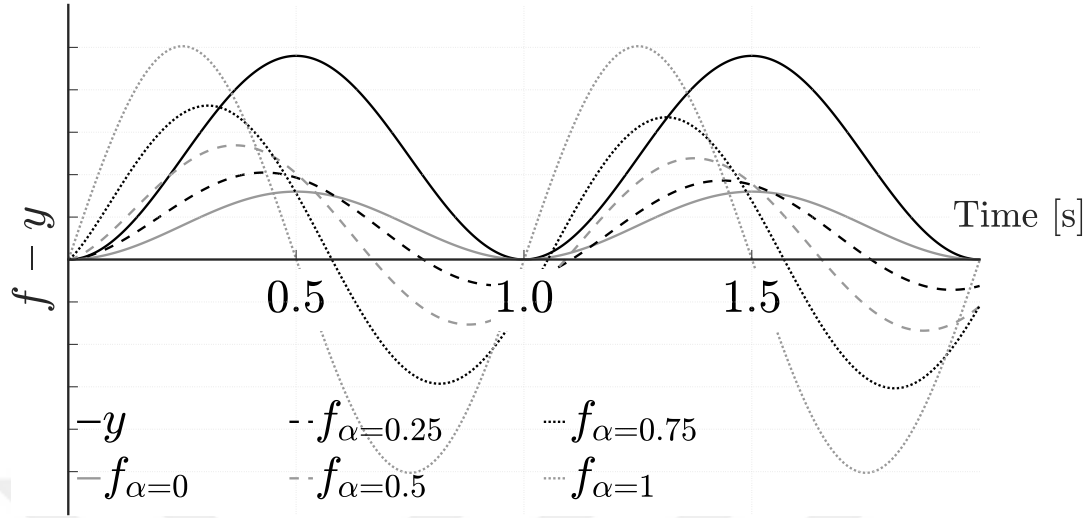


Figure 2.4: Phase difference between the displacement y (the black solid line) and force f for different derivative orders α . For $\alpha = 0$, the y and f (grey solid line) are in phase, and as $\alpha \rightarrow 1$, phase difference increases. Note that for $0 < \alpha < 1$, the peaks of the sinusoidal curves also decrease over time due to relaxation. For $\alpha = 0$, no relaxation occurs; whereas for $\alpha = 1$, relaxation is instantaneous.

$$f(t) = \chi \underbrace{Y e^{i\omega t}}_{y(t)}, \quad (2.14)$$

which is equivalent to a spring equation, whereas, for $\alpha = 1$, the same equation reduces to

$$f(t) = \chi \underbrace{\omega Y e^{i\omega t + \pi/2}}_{\dot{y}(t)}. \quad (2.15)$$

This is equivalent to a dashpot equation where $\dot{y}(t)$ is the time derivative of the input displacement $y(t)$, i.e., the velocity of the input excitation.

2.2.1 Advantages of Springpot Element

One important advantage of the use of springpot element for psychophysical investigation is the straightforward parameterization of the phase difference and the magnitude of the force response. In the case of a springpot element, the phase difference depends solely on the derivative order α . On the other hand, in the case

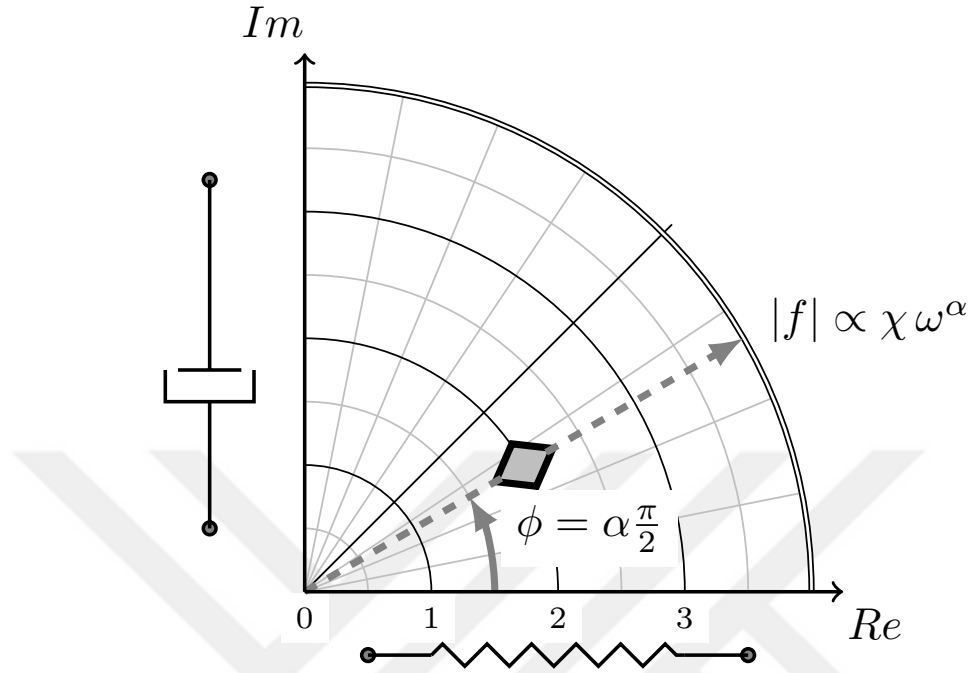


Figure 2.5: Springpot element in polar coordinates. The angular coordinate describes the phase difference ϕ between force and displacement, while the radial coordinate describes the magnitude of the force response. For purely real (imaginary) coordinates, a springpot is equivalent to a spring (damper).

of a Maxwell arm, the phase difference is affected not only by the time constant, which is the ratio of damping and stiffness coefficients, but also the excitation frequency ω (see Eq. 2.7 and Fig. 2.3). Furthermore, in the viscoelastic spectrum, the angular position of a Maxwell arm is defined by its time constant that can take values ranging from 0 to ∞ , whereas the angular position of a springpot is defined by its derivative order α that can take values from 0 to 1. Therefore, the springpot element helps us to define the viscoelastic spectrum more intuitively.

Similarly, while the force response magnitude of a Maxwell arm involves nonlinear combination of the model parameters, the force response magnitude of a springpot element depends on the parameter χ and the power α of the excitation frequency. Depending on the value of α , χ can be equivalent to a stiffness ($\alpha = 0$) or a damping ($\alpha = 1$) coefficient.

2.3 Perceptual Terms for Viscoelasticity

The force-displacement relationship of the springpot element can be described by perceptual terms using the proper choice of adjectives relating to the perceived physical quantities. The springpot model is defined by two parameters: χ and α . Hence, the force response of a springpot element can be described by two perceptual terms relating to the model parameters.

The parameter χ , which scales the force response, is defined as the “firmness” by Scott-Blair and Coppen [44]. The intermediary nature of the word firmness serves the purpose of defining a range between soft and hard materials. According to the online version of the Cambridge Dictionary, the literal definition of firmness is “the quality of not being soft, but not completely hard” [61].

The remaining parameter α determines the angular coordinate of a springpot element on the polar coordinates (see Fig. 2.5). The physical meaning of the angular coordinate can be interpreted in several different ways: phase difference between input and output, the ratio of energy storage to dissipation capacity, or a transition from being a spring to a dashpot. In either way, this parameter defines a springpot’s ability to rebound to its original shape. This quality is described by “springiness” [37, 62], [63] or “bounciness” [62, 64]. We avoid the term “springiness” which is derived from spring and sometimes used synonymously with elasticity [65, 64], and instead, choose to use the term “bounciness” in our study.

Chapter 3

DISCRIMINATION OF BOUNCINESS

In this experiment, we measured the JND for the bounciness and show how bounciness discrimination performance changes along the angular coordinate of the viscoelastic spectrum.

3.1 Experimental Design

We designed a signal detection experiment using one-interval two-alternative forced-choice (1I-2AFC) method. In our experiment, we assumed that the perceived bounciness (i.e., the perception of phase difference) comes from a normally distributed probability density function along the viscoelastic spectrum.

We divided the viscoelastic spectrum of a springpot element into four equal ranges in terms of the phase difference ϕ between force response and displacement input. The beginning of each range corresponds to a reference stimulus S_1 with $\phi = \alpha_1 \pi/2$, where the fractional derivative order $\alpha_1 \in \{0, 0.25, 0.5, 0.75\}$. For each S_1 , we had three comparison stimuli S_2 with derivative orders α_2 with equally spaced increments of $\Delta\alpha$, i.e., $\alpha_2 = \alpha_1 + \Delta\alpha \times n$, where $n \in \{1, 2, 3\}$ (see Fig. 3.1).

In total, there were 12 pairs of stimuli (four references $\times$ three comparisons per reference). Each reference-comparison pair was tested in a different experimental session, and each session was conducted on a different day. This experiment took 12 days to complete for each participant. To randomize the testing order of the four ranges, i.e., reference stimuli, and the three comparison stimuli nested within the ranges, we used the four-by-four Graeco-Latin square design matrix given in Table 3.1 [66]. Here, the Greek letters denote different ranges: $\beta := \alpha_1 = 0$, $\gamma := \alpha_1 = 0.25$, $\zeta := \alpha_1 = 0.5$, and $\theta := \alpha_1 = 0.75$. The Latin letters denote different permutations of the 3 comparison stimuli: $A = \begin{pmatrix} 1 & 2 & 3 \\ 1 & 2 & 3 \end{pmatrix}$; $B = \begin{pmatrix} 1 & 2 & 3 \\ 2 & 3 & 1 \end{pmatrix}$;

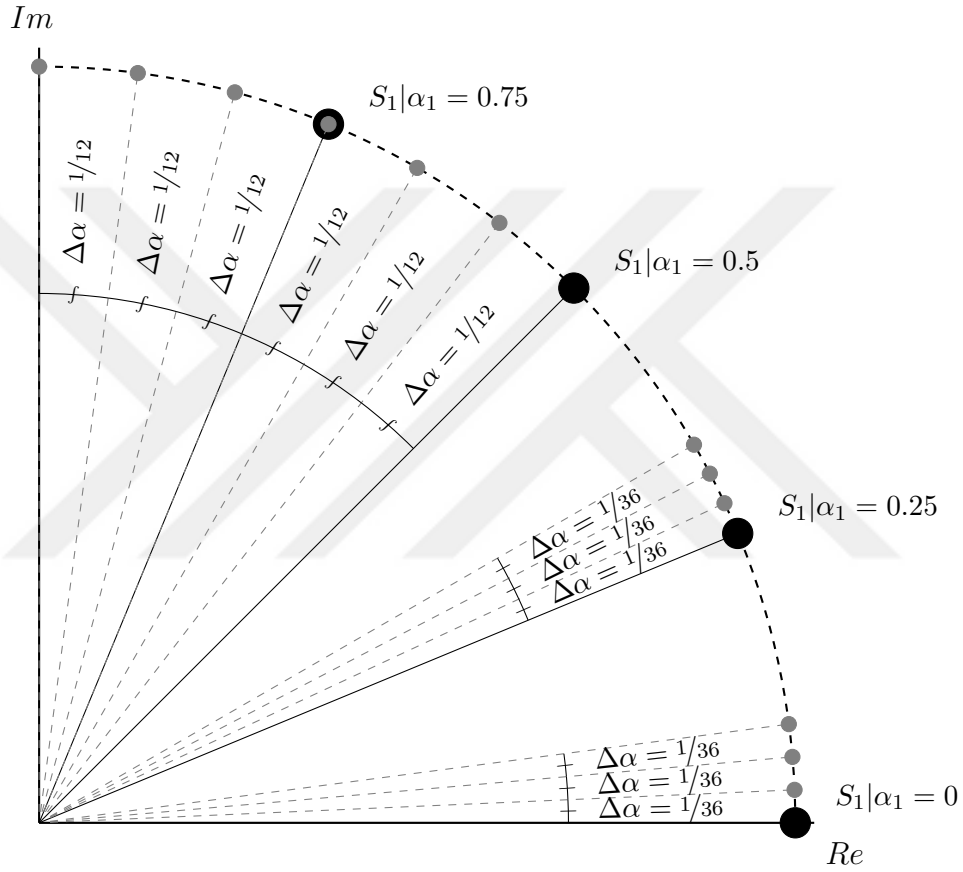


Figure 3.1: The viscoelastic spectrum, and the tested ranges for bounciness experiment. The beginning of each range corresponds to a reference stimulus S_1 (black circle), and three comparison stimuli S_2 (grey circles) were chosen within each range. The reference and the comparison stimuli differ in the angular coordinates by $\Delta\alpha$. The magnitudes of the force response, i.e., the radial coordinates, were equated according to Eq. (3.2).

Table 3.1: Graeco-Latin Square Design matrix

		Participant No.			
		1	2	3	4
Order	1	βA	γD	ζB	θC
	2	γB	βC	θA	ζD
	3	ζC	θB	βD	γA
	4	θD	ζA	γC	βB

$$C = \begin{pmatrix} 1 & 2 & 3 \\ 3 & 1 & 2 \end{pmatrix}; D = \begin{pmatrix} 1 & 2 & 3 \\ 3 & 2 & 1 \end{pmatrix}.$$

3.2 Apparatus

Visual and haptic stimuli were rendered using a 21.5" computer screen and a Phantom Premium 1.0 device with the thimble-gimbal attached, respectively (see Fig. 3.2). To map the workspace of the haptic device to the graphics scene, we used a homogeneous transformation consisting of a pure translation (no rotation). In addition, we projected the graphics scene to the visual display using the orthogonal projection method to avoid any perspective effects.

3.3 Stimuli

All stimuli were springpot elements with a force response along the y -axis only. The graphical representation of the stimuli consisted of a yellow-colored plate and a cyan-colored helical spring under the plate (see Fig. 3.2). The force response of a stimulus was computed numerically in real-time using our “modified” Grunwald-Letnikov method, as explained in the next section.

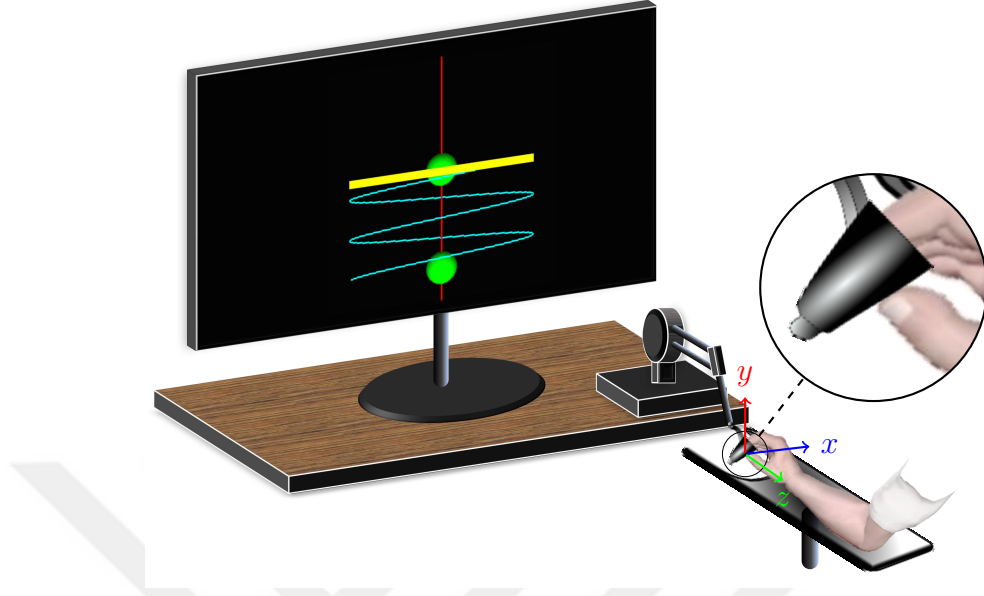


Figure 3.2: Experimental setup and the graphical representation of a stimulus. The participants gripped the end-effector of the haptic device in precision grip pose by placing the index finger inside the thimble-gimble and supporting it with the thumb (see the magnified part in the circle). The stimulus surface and body are represented by the yellow-colored rectangle and the cyan-colored helix, respectively. Note that in the familiarization phase of the experiment, the surface of the reference and comparison stimuli are color-coded: red for reference and green for comparison (see Section 3.1). Additionally, the movement direction is shown with the red line along the y axis which restricts the participant's movement with a snap constraint, and the green spheres denote the starting and ending positions of the movement.

3.4 Numerical Implementation of Fractional Derivative

In Section 2.2, we have presented the solution of the fractional derivatives only for sinusoidal inputs. However, during the experiments, excitation of the viscoelastic stimuli would not necessarily exactly match to the sinusoidal reference that we asked the participants to follow. Hence, we implemented the Grunwald-Letnikov derivative, which is a discrete time formulation of the fractional derivative [57, 58, 67]. Grunwald-Letnikov derivative of an arbitrary function of time $g(t)$ is

$${}_a D_t^\alpha g(t) \approx \frac{1}{\Delta t^\alpha} \sum_{j=0}^{\lfloor (t-a)/(\Delta t) \rfloor} w_j^{(\alpha)} g(t - j \Delta t), \quad (3.1)$$

where

$$w_j^{(\alpha)} = (-1)^{\alpha} \binom{\alpha}{j},$$

and $\lfloor \cdot \rfloor$ is the floor operator. This fractional derivative of order α is defined from time a to time t with a discrete time step of Δt . A close inspection of Eq. (3.1) reveals that the fractional derivatives depend on the entire history of the function of interest. As time progresses, this method becomes both computationally and memory-wise intensive. Since haptic rendering requires real-time force updates at a rate of 1 kHz, the computation of fractional order derivatives with the standard Grunwald–Letnikov derivative can cause problems for our purposes.

A common method to remedy this problem is the short-term memory (STM) principle [68]. According to STM principle, one can set a limit to the number of data points from the history of the differentiated function, and rely on the most recent data which have higher weights. However, we observed that stress–relaxation behavior was not captured well using this approach. Instead, we chose to use a limited number of data points by successive downsampling during the computation of the fractional order derivatives. With the number of data points limited to $N = 100$, the discrete sampling rate at the k^{th} iteration was reduced to $\lceil k/N \rceil$ instead of performing the computation with the whole deformation history. Note that the computations were still performed at the haptic rate of 1 kHz, but the deformation history was sampled at a reduced rate, i.e., Δt was updated such that the number of samples from the history were limited to N . For example, at exploration time $t = 1600$ ms, we sampled the displacement data with a period of $\Delta t = 16$ ms ($\lceil 1600/100 \rceil$), and at the next iteration, i.e., $t = 1601$ ms, we sampled the displacement data with a period of $\Delta t = 17$ ms ($\lceil 1601/100 \rceil$). This meant that for a 3 s exploration, Equation 3.1 was computed with a sampling frequency of 30 Hz, which is well above the human hand movement rate. So, our approach not only captures the stress-relaxation behavior better than an implementation of the Grunwald–Letnikov derivative using short-term memory principle but also lowpass filters the numerical derivative, which results in a smoother derivative function (see Fig. 3.3).

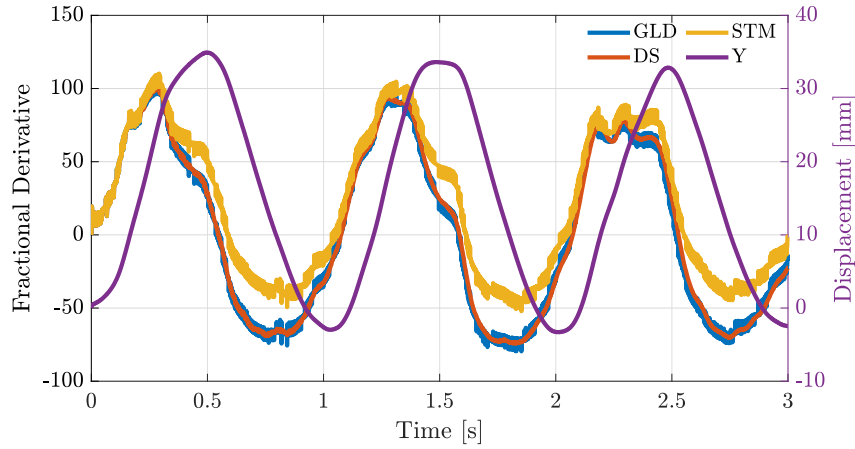


Figure 3.3: The standard Grunwald–Letnikov derivative (GLD) with 3000 data points, downsampled (DS) version with 100 data points, and short-term memory (STM) version with 100 data points. The input displacement (Y) is also shown in purple color with its axis label given on the right.

3.5 Experimental Procedure

Before each session, the participants were asked to sit in front of the computer screen and put their forearms on the adjustable arm rest. They put the index finger of their dominant hand inside the thimble of the haptic device. To avoid slippage, they were instructed to support the thimble with their thumb in a precision grip pose. The haptic interface point was displayed by a blue cursor on the screen.

While exploring the stimuli, the participants were required to follow the procedure depicted in Fig. 3.4. First, they were asked to touch the red-colored snap line (Fig. 3.4a) to constrain their movement in the normal direction. A light guiding force was applied to bring them back to the snap line when their finger deviated from it. Then, they were asked to touch the red colored cursor on the surface of the plate, whose color turned to yellow upon contact (Fig. 3.4b) and then to green after 0.5 s (Fig. 3.4c) to signal that the trial is ready to go. Finally, they were asked to follow the magenta-colored reference cursor, which appears on the screen once the pre-set threshold velocity was exceeded by the participant (Fig. 3.4d). The trajectory of the reference cursor was a sinusoid with a frequency $\omega = 2\pi \text{ rad/s}$. The participants followed this reference cursor and explored each stimulus for 3 s.

During the exploration, the participants always stayed in contact with the stimulus surface, which behaves like a rigid body in that it moves only along the participants' exploration direction and does not deform locally.

We imposed a fixed excitation frequency to attenuate the cues due to the difference in force magnitude $|f| \propto \chi \omega^\alpha$ (see 2.12). We adjusted the firmness coefficient χ_2 of the comparison stimuli S_2 with respect to the firmness coefficient χ_1 of the reference stimuli S_1 as:

$$\chi_2 = \chi_1 \frac{\omega^{\alpha_1}}{\omega^{\alpha_2}}, \quad (3.2)$$

so that the force magnitudes for S_1 and S_2 were equated.

Before the beginning of the experiment, the term “bounciness” was explained to the participants, and they were informed about the nature of the two stimuli in that the reference stimulus S_1 felt bouncier than the comparison stimulus S_2 . Every session started with a familiarization phase where participants were asked to explore the two stimuli until they were able to differentiate them. During this phase, the participants could select and explore S_1 (colored as red) or S_2 (colored as green) without any limitations on the number of explorations per stimulus or the number of switches between the two stimuli. However, simultaneous exploration was not possible. The familiarization phase also served as a training for the exploration procedure. Upon completion of the familiarization phase for the first session, the participants were asked whether they could comprehend the difference between the stimuli and felt comfortable with the term “bounciness” to describe the perception of the stimuli. All participants considered “bounciness” as appropriate for describing the stimuli.

The testing phase started right after the familiarization phase, and each session consisted of 120 trials. In each trial of the testing phase, either S_1 or S_2 was randomly presented to the participants with a neutral color (yellow). The participants' task was to explore the stimulus and respond according to the color codes learned in the familiarization phase: red for bouncier and green for otherwise. Halfway through

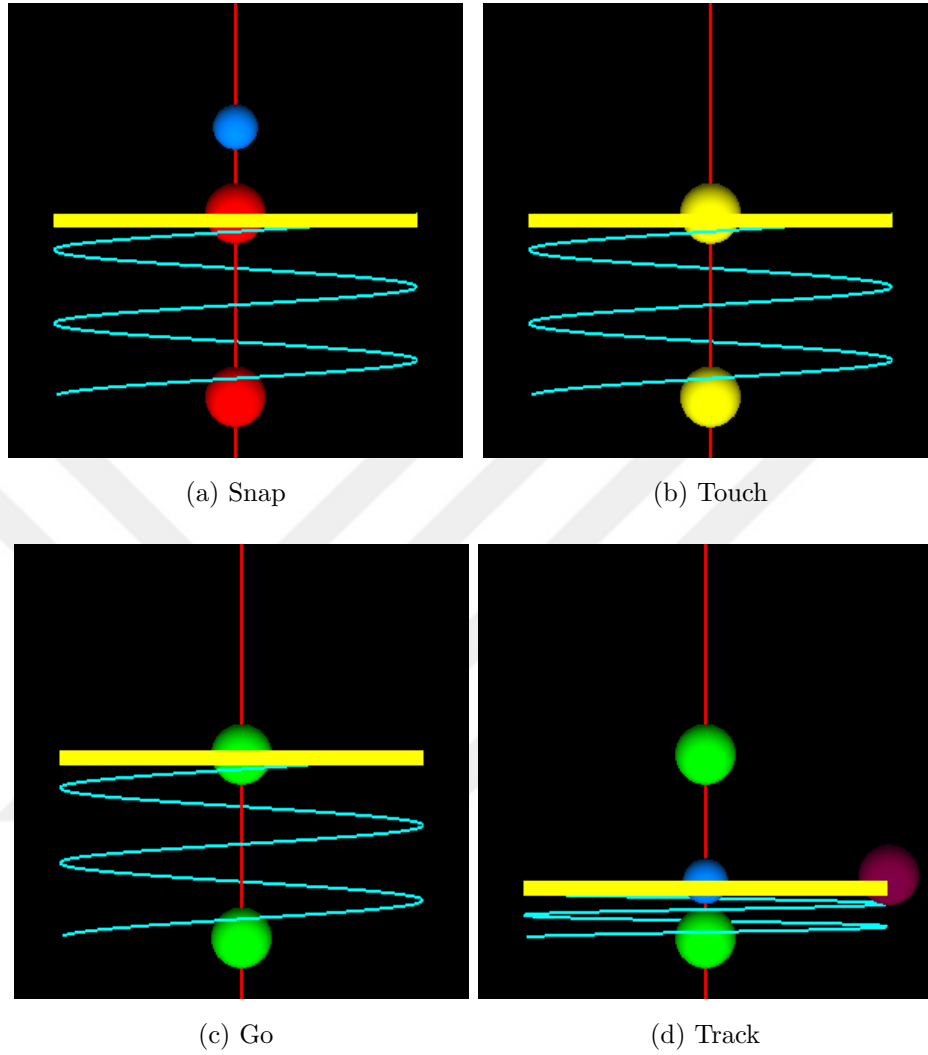


Figure 3.4: Procedural steps of the experiments

the testing phase, a short break was given to avoid fatigue. The first 10 trials at the beginning of the two halves of the testing phase were discarded as training trials. All data analyses were performed on the remaining 100 trials.

3.6 Participants

Four healthy participants (all engineering graduate students, non-native fluent English speakers, right-handed, one female and three males with an average age of 30.25 ± 8.2 years old) performed the experiments. Each participant spent a total

of approximately 6 h (12 conditions $\times$ 30 min, tested on 12 different days) for the whole experiment. We chose to perform our experiments with well-trained and dedicated subjects for a relatively long testing period. The study was approved by the Koc University Human Ethics Committee, and the participants gave their signed informed consent.

3.7 Analysis of Perceptual Data

Based on the participants' responses, we formed two-by-two response matrices for each reference-comparison pair (see Table 3.2). Then, we computed the sensitivity index d' as

$$d' = z(H) - z(FA),$$

where $z(\cdot)$ is the inverse of the cumulative distribution function of the standard Gaussian distribution, and H and FA are the hit and false alarm rates obtained from the response matrix, respectively (see Table 3.2). To avoid the z function to take a value of $\pm\infty$, we added an incorrect response to the replies when necessary, i.e., when $H = 1$ or $FA = 0$.

According to the signal detection theory, for a reference value of α_1 , the ratio $\delta' = d'/\Delta\alpha$ should have a constant value that depends on the participant's sensitivity to the difference between the reference and comparison stimuli. Thus, for each reference stimulus, the three pairs of $(\Delta\alpha \times n | n \in \{1, 2, 3\}, d')$ should be on a straight line with the slope δ' and an intercept at 0. Due to the noise in measurements, one might not get a constant δ' , but instead can calculate the mean of the slopes $\bar{\delta}'$ using the three δ' values. The reciprocal of the mean slope $\bar{\delta}'$, which corresponds to the $\Delta\alpha$ value when $d' = 1$, is the just-noticeable difference: $\text{JND} = 1/\bar{\delta}'$. Also, the mean JND for the population was found from the mean of the slopes across all participants: $\text{JND} = 1/(\mu(\bar{\delta}'))$ [69, 48].

Table 3.2: Stimulus-Response Matrix

		Response	
		R_1	R_2
Stimulus	S_1	CR (Correct reject)	FA (False alarm)
	S_2	M (Miss)	H (Hit)

3.8 Analysis of Physical Interaction Data

We also investigated the relationship between the responses of the participants and several different parameters derived from the recorded physical interaction data, i.e., force and displacement signals. Since we use the term bounciness to describe the sensation related to the fractional derivative order α , we expected the participants to respond in accordance with the phase difference between the reaction force and displacement. It is also conceivable that, while providing their responses, the participants might have relied on cues other than the phase difference. Previously, Tan et al. used the sensitivity index d' as an indicator of the role that mechanical-work cues played when compliance JND was measured with a roving displacement task [48, 7]. They found the d' values to be low and variable for some conditions when assuming that the participants were responding to compliance cues. When the same data were reprocessed assuming that the participants were using mechanical-work cues, a higher and most stable d' value emerged. This result was taken as evidence that the participants may have relied more on mechanical-work cues than compliance cues when performing compliance discrimination [48].

In this work, we adopted the approach in [48] due to the similar experimental methods. We estimated the average (physical) distances $\hat{d}'$ between the reference and comparison stimuli for several candidate parameters and then normalized these values with respect to the corresponding (perceptual) d' values. For clarity, we used the term “sensitivity index” with the perceptual data and “normalized distance” with physical parameters even though they were computed in similar ways.

The physical interaction parameters were chosen based on our viscoelastic model and also the earlier studies where similar analyses were performed [48, 10, 11]. The selected parameters were categorized into frequency-domain parameters:

- k_S : Real (storage related) part of complex stiffness
- k_L : Imaginary (loss related) part of complex stiffness
- η : Loss factor $(k_S)/(k_L)$
- b_L : Imaginary part of force over the magnitude of velocity (Imaginary part of the mechanical impedance)
- ϕ : Phase difference between force and displacement,

and time-domain parameters:

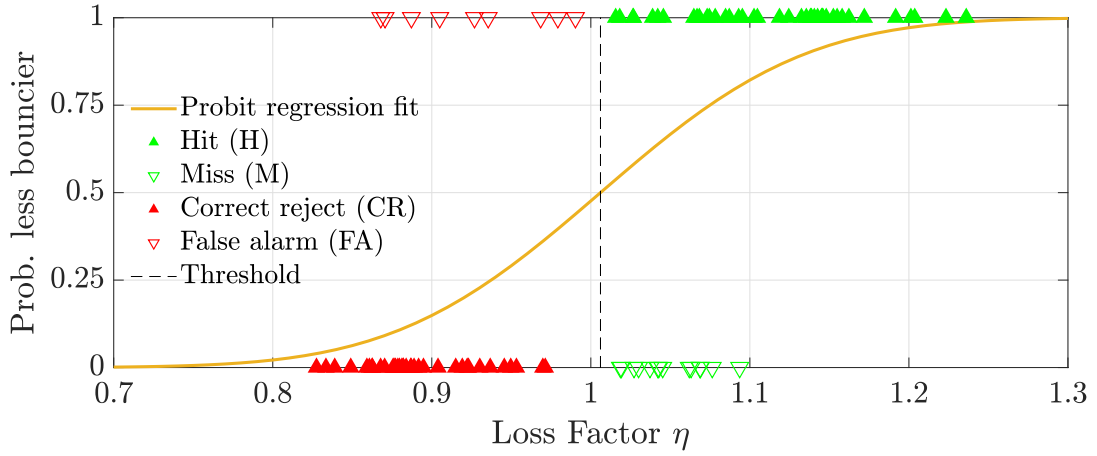
- W : Positive part of the mechanical-work done by the participant
- f_{max} : Maximum of the force peaks
- $\bar{f}$: Mean of the force peaks
- y_{max} : Maximum of the displacement peaks
- $\bar{y}$: : Mean of the displacement peaks
- v_{max} : Maximum of the velocity peaks
- $\bar{v}$: Mean of the velocity peaks
- k_I : Initial local stiffness (first force peak divided by the corresponding indentation depth)
- k_{fit} : Slope of the zero-intercept line fitted to force as a function of displacement (force over position, FOP, as in [10])

- b_{fit} : Slope of the zero-intercept line fitted to force as a function of velocity [10].

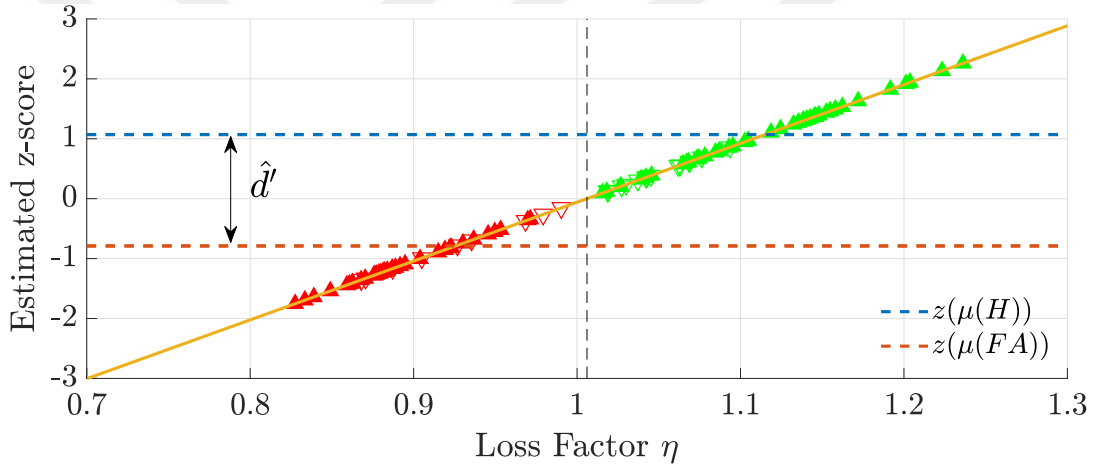
The frequency-domain parameters were computed based on the most dominant frequency components of the relevant data. The frequency components of the data and their magnitudes were identified using the Discrete Fourier Transform (DFT) functionality of Matlab. Following the DFT, we adjusted the phase of each signal using a rotation matrix that took the displacement data as the reference, i.e., the transformed displacement data always had a phase of 0 rad.

Some of the selected physical interaction parameters depend on the programmed material properties of the stimuli and the motion of the participants who were instructed to follow a visual cursor. Since our exploration procedure was prescribed, one might think that the values of these parameters can be pre-computed and then the normalized distances can be easily computed using the methods from the previous section. However, the participants' performance in following the visual cursor determines the actual values of the interaction parameters derived from the recorded force, displacement, and velocity data. So, the physical interaction parameters do not match the expected ones perfectly. Because of these deviations or the parameter's nature, the recorded physical interaction data do not necessarily form clearly distinct clusters. Therefore, we cannot use the methods described in the previous section to compute normalized distances based on the physical interaction data.

To investigate the relation between the physical data and the responses of the participants, we fitted generalized linear models (GLM) using a Bernoulli distribution and a probit link function (see Fig. 3.5a). The fitted linear models transform the participants' responses (0 or 1) to z-scores as functions of physical interaction parameters (see Fig. 3.5b). Using these linear models, we classified the stimuli experienced by the participants as noise or signal depending on the sign of the estimated z-scores. If the estimated z score of a trial is positive (negative), then this trial is classified as a signal (noise). Then, we can estimate a parameter specific normalized distance $\hat{d}'$ between the means of the physical parameters in z coordinates.



(a) Normal cumulative function fit to data



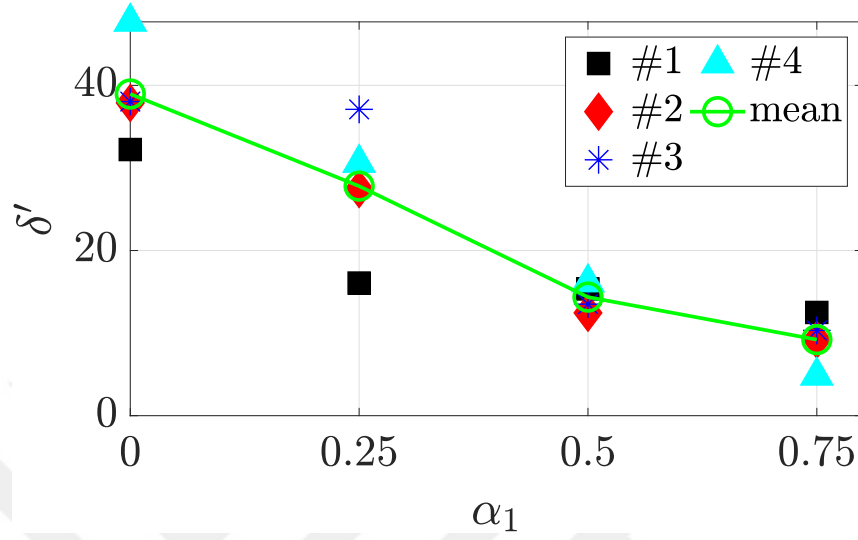
(b) Data mapped to the fitted line

Figure 3.5: An example showing the curve fit to loss factor data versus participant responses using probit function. (a) The triangles denote the data points, i.e., the loss factor data of the trials vs the participant responses (1: Less bouncy, 0: Bouncer). (b) The data points are mapped to the fitted line as z scores. Based on the fit, the green data points are classified as less bouncy stimuli and the red ones as the bouncer.

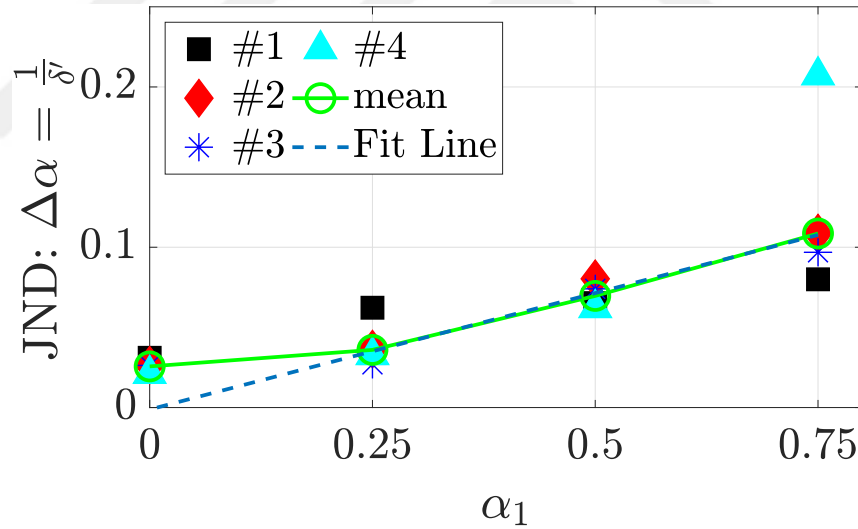
3.9 Results

The slopes $\bar{\delta}'$ and their reciprocals (JNDs) are plotted in Fig. 3.6a and Fig. 3.6b, respectively. A linear relationship was observed between the mean JND values of the whole group and the reference $\alpha_1 > 0$, and a linear regression model was fitted using Matlab (see Fig. 3.6b). The fitted line has a slope of 0.1452 ($t(1) = 20.0065, p =$

0.0318) and an intercept of -0.0012 ($t(1) = -0.3088, p = 0.8093$) with $R^2 = 0.998$.



(a) Slopes δ'



(b) JNDs

Figure 3.6: The slope (a) and JND (b) results of each participant and the group means for all references. In 3.6b, a blue dashed line is fitted to the mean JNDs of references for $\alpha_1 > 0$

The above results were computed with respect to changes in the derivative order α which results in a change in phase difference ϕ . Nevertheless, the participants might have based their decisions on the differences in other physical parameters

arising from the interaction with the stimuli. To investigate the possible role of these parameters in their perception, we computed the ratio of normalized distance over the sensitivity index ($\hat{d}'/d'$) for several physical interaction parameters and averaged the resulting values across the increments and the participants for all references. In Fig. 3.7, for each reference, we report the three parameters having the largest ratios between signal and noise.

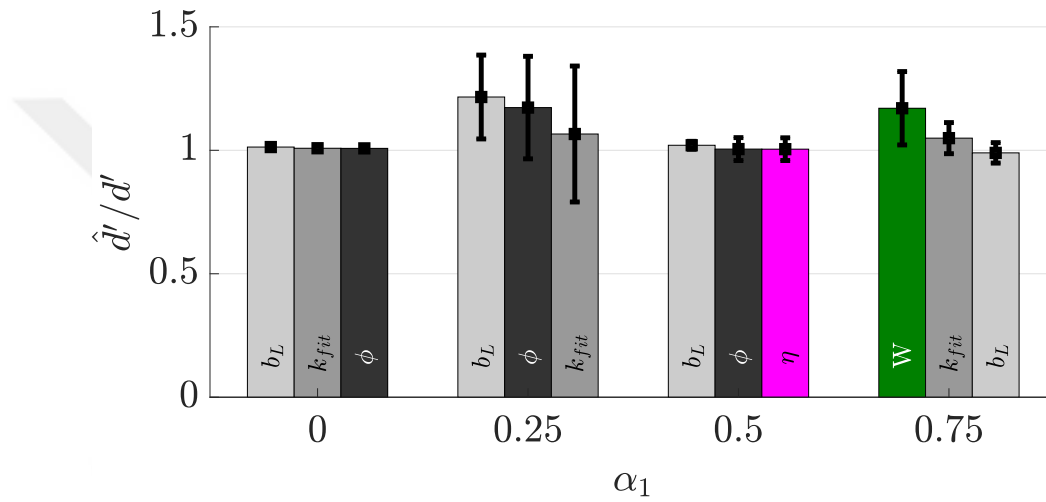


Figure 3.7: The ratios of estimated $\hat{d}'$ of the physical parameters to d' in Experiment I. The results are the mean values averaged over the participants and the 3 $\Delta\alpha$ increments, and the errorbars show the standard errors of the means. For each reference, only the three parameters with the largest ratios are reported.

We observed that imaginary part of the mechanical impedance b_L was among the top three parameters for all references. Also, the ratio for the phase difference ϕ was at least 0.97 for all references. In contrast, for $\alpha = 0.75$, mechanical-work W had the highest ratio.

Chapter 4

DISCRIMINATION OF FIRMNESS

Another perceptual characteristic of a springpot element is its firmness. Assuming that bounciness perception does not change for a fixed value of α , we can characterize firmness solely by χ . In this experiment, we measured the JNDs for the firmness of springpot elements that have a fixed location along the viscoelastic spectrum in terms of phase difference, i.e., the different reference stimuli were characterized by their α values while the comparison stimuli differed from the references only in χ (see Fig. 4.1). With this experiment, we aimed to fill the gap in the literature between the stiffness (firmness when $\alpha = 0$) and damping/viscosity (firmness when $\alpha = 1$) JND values.

4.1 Methods

We utilized the same apparatus, experimental design, and data analysis methods as in the bounciness experiment. However, one major difference was in the exploration procedure. In the current case, a change in χ does not affect the angular coordinate of a springpot element along the viscoelastic spectrum. Therefore, unlike in the previous case, we do not have to impose a specific movement frequency on the participants to attenuate extra cues. This freedom results in a visuo-haptic interaction that is much closer to our daily life experience.

Another difference is the inclusion of $\alpha = 1$, which increased the number of reference stimuli S_1 to five with the derivative orders $\alpha_1 \in \{0, 0.25, 0.5, 0.75, 1\}$. We again had three comparison stimuli S_2 with firmness coefficients $\chi_2 = \chi_1 + \Delta\chi$ (see Fig. 4.1).

The firmness coefficient of the reference stimuli was normalized as

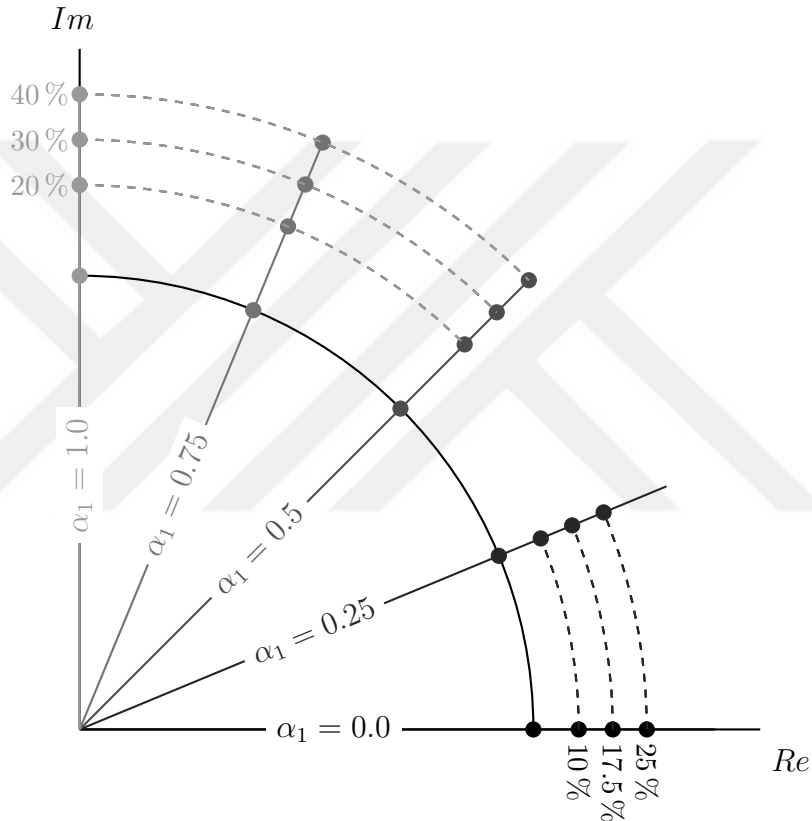


Figure 4.1: The viscoelastic spectrum, and the tested ranges for firmness experiment. The beginning of each range corresponds to a reference stimulus S_1 (large circles), and three comparison stimuli S_2 (smaller circles) are chosen within each range. The comparison stimuli have the same α values as their respective reference stimuli. The reference and the comparison stimuli differ in the radial coordinates by $\Delta\chi$. The levels of $\Delta\chi$ are shown with the grey arcs.

$$\chi_1 = \chi_0 \frac{\omega^{\alpha_0}}{\omega^{\alpha_1}} = \frac{\chi_0}{\omega^{\alpha_1}}, \quad (4.1)$$

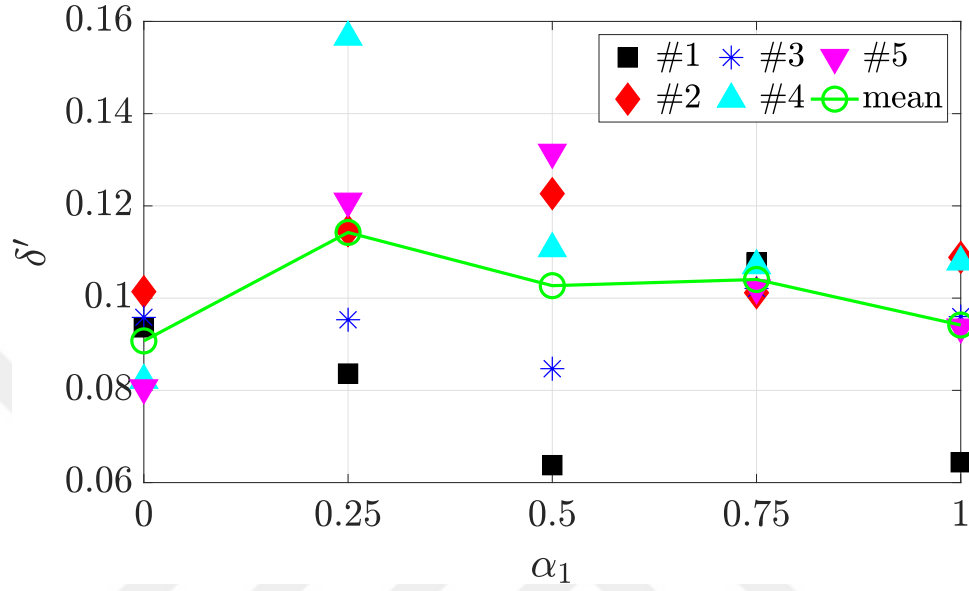
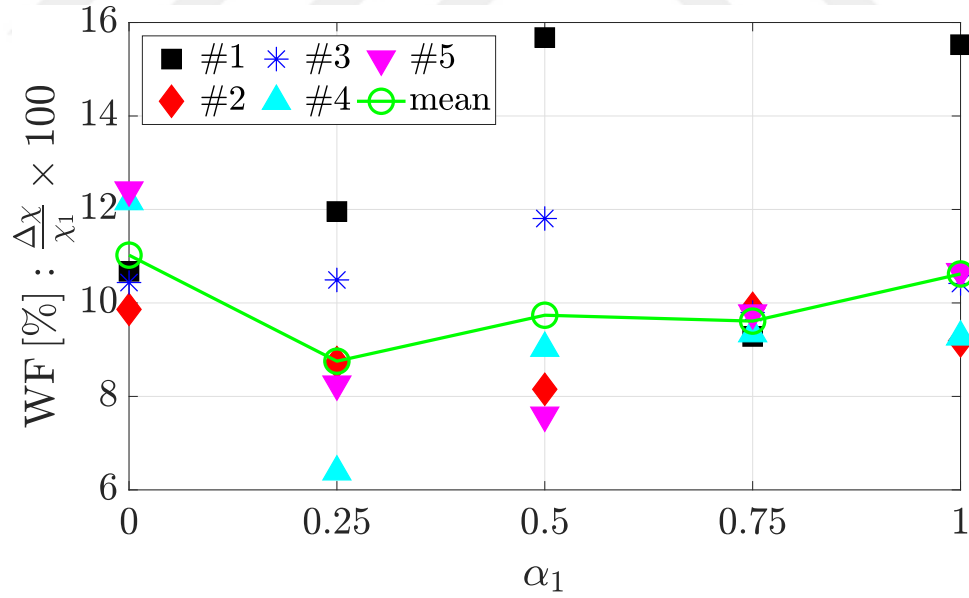
where $\omega = 2\pi \text{ rad/s}$ was chosen as a reference excitation frequency, $\alpha_0 = 0$, $\chi_0 = 0.3 \text{ N/mm}$, and the unit of χ_1 is $\text{N} \cdot \text{s}^{\alpha_1}/\text{mm}$. This normalization was performed for achieving comparable force responses across the reference stimuli.

The above changes increased the number of pairs (testing conditions) to 15 (5 references $\times$ 3 comparisons). Consequently, the total experimentation time increased to approximately 7.5 h (15 conditions $\times$ 30 min) per participant. The size of the Graeco-Latin Square design matrix was increased to 5×5 to accommodate the increase in the number of references. Therefore, we needed to have an equal number of references, permutations of increments, and participants. So, we included an additional permutation $E = \begin{pmatrix} 1 & 2 & 3 \\ 2 & 1 & 3 \end{pmatrix}$ and recruited five new participants (all engineering graduate students, right handed, and males with an average age of 32.2 ± 5 years old).

4.2 Results

The slopes of the lines fitted to the d' and the JND values as a percentage of the reference χ_1 values, i.e., Weber fractions (WFs), are depicted in Fig. 4.2a and 4.2b, respectively. The slopes and hence the JND values do not follow a clear trend with respect to α_1 unlike in the first experiment. A repeated measures ANOVA in SPSS showed no significant effect of α_1 ($F(4, 16)$, $p = 0.3981$), i.e., level of bounciness on firmness WF. Also, no significant change has been observed for the pairwise comparisons with respect to α_1 levels (all $p > 0.5$).

As in the previous experiment, we computed the ratios of the estimated $\hat{d}'$ of the chosen physical interaction parameters to the d' of $\Delta\chi$ and averaged these ratios across participants. In Fig. 4.3, the averaged ratios of the physical parameters with the three largest values are reported. According to this plot, frequency-domain parameters (k_S for $\alpha = 0, 0.25$ and 0.5 , and b_L for $\alpha = 0.75$ and 1), and regression parameters (k_{fit} for $\alpha = 0, 0.25$ and 0.5 ; and b_{fit} for $\alpha = 0.5, 0.75$ and 1) appeared to have the largest ratios between signal and noise stimuli.

(a) Slopes δ' 

(b) JNDs in Weber fractions

Figure 4.2: (a) The slope and (b) JND in Weber fractions results of each participant and the group means for all references.

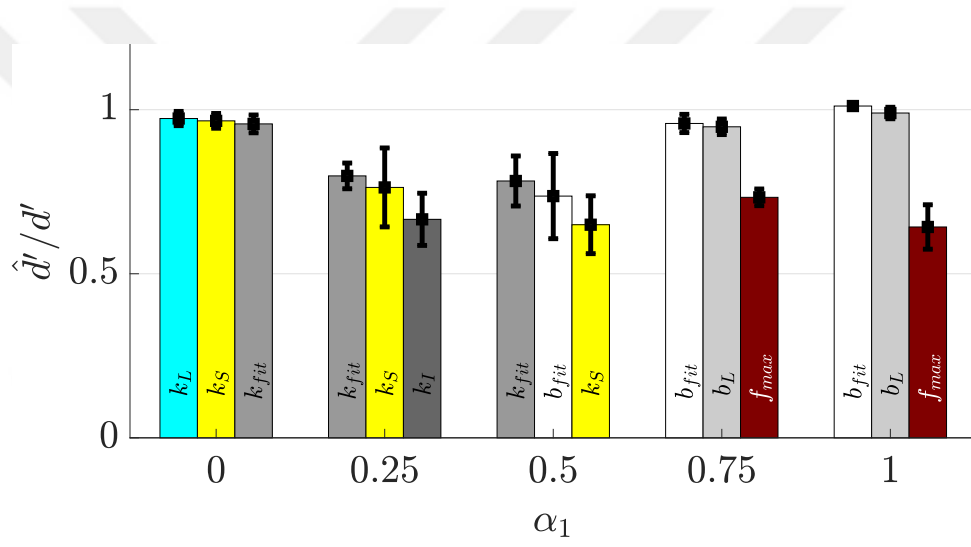


Figure 4.3: The ratios of estimated $\hat{d}'$ of the physical parameters to d' in Experiment II. The results are the mean values averaged over the participants, and the errorbars show the standard errors of the means. For each reference, only the three parameters with the largest ratios are reported.

Chapter 5

DISCUSSION

In this study, we have considered viscoelasticity as a mechanical response displaying a spectrum of behaviors varying from pure elastic (spring) to pure viscous (damper). We also adopted the “bounciness” and “firmness” from the psychorheology field as the perceptual terms describing the sensations associated with viscoelastic materials. At various locations on the viscoelastic spectrum, we measured the bounciness and firmness JNDs via two psychophysical experiments.

5.1 Bounciness Experiments

In the first experiment, bounciness JND was found to be monotonically increasing with the α parameter of the springpot element. For references with α_1 values larger than zero, the mean JND values followed a linear trend with its intercept being approximately zero (see Fig. 3.6b). This trend is reminiscent of Weber’s law which states that the JND of a physical quantity is a constant fraction of the reference value. We should also note that transforming our modeling-basis might result in a change in the measured “normalized distances” between the stimuli categorized as signal and noise. If the participants made their decisions based on a specific cue, the estimated normalized distance for this cue should be larger than the other estimated distances. This argument acknowledges the fact that although the experimenter has clear definitions of signal and noise based on the model, the participants might have their own definitions based on the sensory cues they rely on. The ratio of the normalized distance to the sensitivity index ($\hat{d}'/d'$) tells more about this argument (see Fig. 3.7). A ratio lower than one ($d' > \hat{d}'$) suggests that the parameter manipulated by the experimenter (α in bounciness and χ in firmness experiments) is more effective than the physical parameter in question on the participants’ perception. Analysis of this ratio revealed that the frequency-domain parameters were the ones

most likely affecting bounciness perception for $\alpha_1 < 0.75$. In this regard, the imaginary part of the mechanical impedance b_L was always in the top three parameters. Except for the last reference, the phase difference ϕ was also among the top three parameters, and its ratio was at least 0.97 for all references. However, for the last reference, mechanical-work W had the highest ratio (see Fig. 3.7 for $\alpha_1 = 0.75$).

Based on these findings, we can explain the increase in JND with respect to α by the energy dissipation characteristics of the stimuli which is related to the imaginary part of the mechanical impedance b_L . In the case of a pure spring, the addition of even some dissipation was easily detected by the participants. However, as α grows, both the reference and comparison stimuli became more dissipative, and the participants might have had difficulty identifying the relatively more dissipative stimulus. Instead, they might have relied on the cues related to effort, i.e., mechanical-work W .

5.2 Firmness Experiments

In the firmness experiment, we did not observe a significant change, let alone an increase in the JND in Weber fractions WFs as the stimuli ranged from a pure spring ($\alpha = 0$) to a pure damper ($\alpha = 1$). Our results show that the firmness WFs does not change when the phase difference is fixed across the stimuli. Nevertheless, this does not necessarily mean that firmness discrimination of viscoelastic stimuli is independent of its viscous component. In fact, a delay in force feedback (a phase difference in force response with respect to displacement input) has already been shown to affect both the perceived stiffness (firmness in our context) and discrimination performance [13, 10, 30]. Nonetheless, as discussed in [70], the WFs reported in earlier studies for viscosity (13.6 % [49]; 34 % [51]) are larger than the WFs reported for elasticity (8 % [48]; 23 % [50]).

In terms of the experimental methods, our study is most similar to [48] and [49] where signal detection experiments were performed with the elastic and viscous stimuli rendered by an electromechanical device. In those studies, compliance and viscosity WFs were $\sim 8\%$ and $\sim 14\%$, respectively. Nevertheless, the WF

values reported in [49] have a high standard deviation among participants. Later in [71], Beauregard and Srinivasan acknowledged this high standard deviation and attributed it to the relatively poor discrimination performance of one of the three participants. Furthermore, they included new data from another viscosity JND experiment with different reference stimuli varying in the damping coefficient for a fixed displacement of 30 mm, and the new results showed the viscosity WF to be below 10 %. So, the viscosity JND reported as Weber fractions in [49] seemed to be overestimated.

We can also compare our results with other haptic perception studies on viscosity JND. In [72], the authors investigated the masking effect of an additional stiffness and inertia on the viscosity JND using a programmable electromechanical device. According to the results of the control condition where no masking was applied, the mean WF was reported as being around 10 %. In contrast, in earlier studies where real liquid test specimens were used, the mean WF values were reported as 30 % in [44] and above 30 % in [9]. However, the results have high inter-condition and inter-participant variations, and the WFs of some participants are considerably below 10 % for some conditions. These large variations also support our choice of using a relatively small number of well-trained participants.

To understand the cues used for discrimination of firmness, we computed the normalized distances based on the parameters obtained from the physical interaction data. We observed that the ratio ($\hat{d}'/d'$) was never greater than one (see Fig. 4.3). This result suggests that the participants took the frequency variations in their movements into account when they responded. Therefore, the manipulated parameter χ provided the relevant cues for the perception of firmness.

5.3 Effect of (Presence of) Visual and (Lack of) Tactile Feedback

Although our experimental methods are similar to those used in [48, 49] and [71], two major differences were present in our experiments: (1) visual feedback was displayed to the participants, and (2) the indentation distance was not fixed. Nevertheless, our JND results in Weber fractions are comparable to those reported in [48] and [71].

While the uncontrolled displacement was expected to yield a lower discrimination performance similar to the roving displacement condition (22 % [7]), the inclusion of the visual cues might have aided the participants in our experiments. It is conceivable that our participants compensated for uncontrolled displacement with the help of visual cues. Earlier studies have already shown that reliability in haptic stiffness estimation can be improved with the integration of visual information [12, 73, 74]. Regardless of the sensory integration process, both Wu et al. [53] and Varadharajan et al. [54] reported statistically significant improvement in stiffness JND when visual feedback was available with proprioceptive force feedback.

On the contrary, visual feedback was found not to improve the sensitivity to the difference between reference and comparison stimuli for real specimens [55]. The authors argued that this result was due to the higher reliability of local surface deformation cues in firmness estimation. Real specimens enable a natural interaction and provide valuable cutaneous information which is limited in the case of electromechanical (EM) devices displaying proprioceptive force feedback. The implications of the reduced cutaneous feedback have been studied by using elastic stimulus with a rigid plate attached to its top surface [75, 8]. In [75], Srinivasan and LaMotte showed that participants could discriminate pairs of rubber stimuli much better than pairs of spring cells with rigid plates on their top surfaces. Bergmann Tiest and Kappers measured the softness (firmness) JNDs using rubber stimuli with and without rigid plates attached to the top and bottom surfaces. They reported an increase in Weber fraction from 15 % to 50 % when the rigid surfaces were used, i.e., when the cutaneous cues were reduced.

On the other hand, real specimens are usually prepared by curing silicone rubber mixtures [15, 8, 17, 31, 76], and it is known that silicone rubber exhibits nonlinear and viscoelastic material behavior [77]. Hence, care should be given to their mechanical characterization. Nevertheless, so far, mechanical characterization of these specimens for the psychophysical investigation of haptic perception has been non-standard, and the recent work by Gerling et al. is an important step towards solving this problem [78]. More specifically, tuning the viscoelastic properties of a silicone

rubber while controlling its material nonlinearities is a challenge that has not been solved yet.

Conversely, our work required simultaneous tuning of viscoelastic material properties related to firmness and bounciness. For this reason, we have chosen to use an EM device in the current work. Despite the reduced cutaneous cues, EM devices can still successfully simulate the scenarios involving tool-object interactions. In that respect, our results, for example, could be useful in the design and evaluation of surgical simulators and teleoperated surgical systems where simplifications in tissue models and reductions in data transfer rates are required to meet real-time constraints.

Chapter 6

CONCLUSION

In summary, we observed that the JND values for bounciness increased as a function of the fractional derivative order, whereas the WF for firmness was almost constant across the five equally spaced points along the angular coordinate. When bounciness and firmness are considered as the terms for assessing the perception of viscoelastic materials, we can successfully explain the participants' decision-making using frequency-domain parameters. The frequency-domain parameters themselves provide cues related to the energy storage and dissipation characteristics of the stimuli. Energy storage and loss characteristics relate to a sense of effort which might have been used as an intuitive indicator for the perception of viscoelastic stimuli.

In our future studies, we will explore the viscoelastic spectrum as a bivariate domain, and investigate the effect of interaction between the parameters χ and α on the perceived bounciness and firmness. We also plan to investigate the roles temporal combination and integration of different sensory signals play in firmness perception.

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